

Optimized diamond photonic molecule for quantum communications

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In recent years, elementary quantum optical structures, called photonic molecules (PMs), have been carefully studied both experimentally and theoretically [1, 2]. These structures are formed from high quality factor solid-state microresonators (MR). These devices may be integrated with single-photon sources that generate and guide photon flows in a system and high-sensitivity detectors that fix the arrival of a photon and, preferably, its polarization [3]. As for the element base for quantum computation, the main effort of scientists is now focused on the search for the optimal geometry of a solid-state photonic chip [4]. Here, we propose the design of three-unit PM optimized to obtain good transport and dissipation properties.

To design PM supporting optical-band frequencies, one uses photon cells with geometric dimensions on the order of a few microns. MRs supporting whispering gallery modes (e.g. microrings) can form quasi-one-dimensional optical structures. We optimize diamond microring parameters calculating the eigenfrequencies and the electrical field distributions of the single microring in a broad range of inner and outer radii as well as thicknesses.

Analytical consideration of the PM-system composed of three MRs is given within the formalism of tight-binding phenomenological Hamiltonian:

$$H = \sum_{k=1}^3 (\omega_k - i\kappa_k) a_k^\dagger a_k - \sum_{k=1}^2 J_{k,k+1} (a_k^\dagger a_{k+1} + a_{k+1}^\dagger a_k), \quad (1)$$

where ω_k is the mode frequency of the k -th MR ($k = 1 - 3$), a_k^\dagger and a_k are creation and annihilation operators of photons, respectively, $J_{k,k+1}$ is a coefficient of photon hopping between the MRs, κ_k is a rate of energy dissipation of the MR mode. Provided that $\omega_k \equiv \omega$ and $J_{k,k+1} \equiv J$ each mode of the single MR splits into three ones of PM with frequencies $\omega_{PM1,3} = \omega \pm \sqrt{2}J$, $\omega_{PM2} = \omega$. The electric field profile of PM for the eigenfrequencies $\omega_{PM1,3}$ has antinodes located along the edge of each ring.

The most common mean for controlling the dynamics of photons in PM is laser. Here we employ the weak laser as a probe for the PM's spectrum. The probability of one-photon excitation of PM due to laser photon injection is calculated in the steady-state regime. The transmission spectrum was represented by a small number of clearly distinguishable peaks. As the MRs approach each other, the splitting of the peaks increases. It is very desirable to have *equal* optical field amplitudes in each resonator for some eigenmode of the PM. Our solution consists in following choice of Hamiltonian parameters: $J_{k,k+1} \equiv J$, $\omega_1 = \omega_3 = \omega$, $\omega_2 = \omega + J$. In this case, the mode with the frequency $\omega_{PM,opt} = \omega - J$ is equally-weighted: $|\Psi_{opt}\rangle = (|1\rangle + |2\rangle + |3\rangle)/\sqrt{3}$.

References

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